

Assessing Alkyl-, Silyl-, and Halo-Substituent Effects on the Electron Affinities of Silyl Radicals

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Neutral anion energy differences for a large class of α -substituted silyl radicals have been computed to determine the effect of alkyl, silyl, and halo substituents on their electron affinities. In particular, we report theoretical predictions of the adiabatic electron affinities (AEAs), vertical electron affinities (VEAs), and vertical detachment energies (VDEs) for a series of methyl-, silyl-, and halo-substituted silyl radical compounds. This work utilizes the carefully calibrated DZP++ basis set (Chatgialaloglu, *C. Chem. Rev.* **2002**, 102, 231), in conjunction with the pure BLYP and OLYP functionals, as well as with the hybrid B3LYP, BHLYP, PBE1PBE, MPW1K, and O3LYP functionals. Bromine has the largest effect in stabilizing the anions, and the BLYP/DZP++ AEA for SiBr₃ is 3.29 eV. The other predicted electron affinities are for SiH₃ (1.37 eV), SiH₂CH₃ (1.09 eV), SiH₂F (1.54 eV), SiH₂Cl (1.94 eV), SiH₂Br (2.05 eV), SiH₂(SiH₃) (1.77 eV), SiH(CH₃)₂ (0.92 eV), SiHF₂ (1.86 eV), SiHCl₂ (2.53 eV), SiHBr₂ (2.67 eV), Si(CH₃)₃ (0.86 eV), SiF₃ (2.66 eV), SiCl₃ (3.21 eV), Si(SiH₃)₃ (2.25 eV), and SiFCIBr (3.13 eV). For the five silyl radicals where experimental data are available, the BLYP functional gives the most accurate determination of AEAs; the average absolute error is 0.04₁ eV, whereas the corresponding errors for the O3LYP, MPW1K, PBE1PBE, B3LYP, OLYP, and BHLYP functionals are 0.05₈, 0.06₀, 0.06₃, 0.08₅, 0.11₅, and 0.15₃ eV, respectively.

I. Introduction

Recently, high concentrations of very reactive silyl radicals have been observed in silane plasmas, clearly demonstrating their importance in chemical vapor deposition (CVD) processes.¹ Indeed, SiH₃ radicals are the dominant reactive species in the production of a variety of high-temperature structural materials, including ceramics.² Silyl radicals also play a central role in the free-radical hydrosilation of poly(phenylsilane); this process has applications in the fields of photoconducting materials, microlithography, and nonlinear optics.³ Although little is currently known about how to employ silyl radicals effectively in organic synthesis, the boom in carbon-based radical-driven synthetic reactions, in conjunction with the ability of silyl radicals to form stereoselective bonds, may provide the driving force needed to popularize these radicals soon.⁴ Silyl radical anions (these are closed-shell systems) are typically employed to generate silicon–carbon bonds,^{5,6} and they act as nucleophiles, analogous to their carbon counterparts. Applications involving these anions include the preparation of acylsilanes in sugar chemistry (via copper catalysis reactions with acyl halides⁴) and the production of bioactive compounds that are synthetically equivalent to amino aldehydes.⁷

Computations have already successfully characterized a large set of atomic and molecular electron affinities,⁸ substituent effects on methyl radicals,⁹ and substituent effects on silyl- and carbanion stability.⁵ In this research, we assess the extent to which silyl and halogen substituents impact the structures, energetics, and electron affinities of silyl radicals. In our continuing efforts⁸ to calibrate DFT functionals for the accurate determination of electron affinities, we have also included here computations with the more recently developed OLYP, O3LYP,

PBE1PBE, and MPW1K functionals in those cases where experimental data is available. In addition, we discuss the inverse relationship between electronegativities and electron affinities of silyl radicals and compare to earlier analogous work on methyl radicals⁹ and silylenes.¹⁰

II. Theoretical Methods

Equilibrium geometries and vibrational frequency analyses for all the radicals and anions in this study were predicted using generalized gradient approximation (GGA) exchange correlation density functionals, as well as two HF/DFT hybrid functionals. The three “standard” functionals are designated BLYP, B3LYP, and BHLYP, respectively, representing the pure DFT exchange functional (B),¹¹ a three-parameter HF/DFT hybrid exchange functional (B3),¹² and a modified half-and-half HF/DFT hybrid method (BH),¹³ all coupled with the dynamical correlation functional of Lee, Yang, and Parr (LYP).¹⁴ In addition, for the five species of silyl radicals where reliable experimental electron affinities are available, we employed the pure OLYP and hybrid O3LYP functionals. These functionals incorporate the novel OPTX local exchange functional of Handy, Cohen, and co-workers^{15,16} the hybrid PBE1PBE^{17,18} functional, which utilizes the one-parameter GGA PBE¹⁹ functional and the MPW1K functional. The latter is constructed from a modified version of the Perdew–Wang gradient-corrected exchange functional (MPW)²⁰ and has been parametrized specifically for kinetics.²¹ All DFT computations were carried out using the GAUSSIAN 94,²² 98,²³ and 03²⁴ suite of programs.

Double- ζ quality basis sets with polarization and diffuse functions (DZP++) were used throughout for geometry optimizations and vibrational frequency analyses. The DZP++ basis sets were constructed from the Huzinaga–Dunning–Hay^{25–27} sets of contracted Gaussian functions. A set of p-type polarization functions for each hydrogen atom was added, and one set

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of five d-type polarization functions was included on each heavy atom. Basis functions for bromine were obtained from the standard Ahlrichs' double- ζ spd set²⁸ with one set of more extended d-like polarization functions [$\alpha_d(\text{Br}) = 0.389$]. These basis sets were further augmented with diffuse functions; each atom received one additional s-type and one additional set of p-type functions. The diffuse orbital exponents were determined in an even-tempered sense according to the prescription set forth by Lee and Schaefer:²⁹

$$\alpha_{\text{diffuse}} = \frac{1}{2} \left(\frac{\alpha_1}{\alpha_2} + \frac{\alpha_2}{\alpha_3} \right) \alpha_1$$

where α_1 , α_2 , and α_3 are the three smallest Gaussian orbital exponents of the s- or p-type primitive functions for a given atom ($\alpha_1 < \alpha_2 < \alpha_3$). The added diffuse functions have orbital exponents $\alpha_s(\text{H}) = 0.04415$, $\alpha_s(\text{C}) = 0.04302$, $\alpha_p(\text{C}) = 0.03629$, $\alpha_s(\text{Si}) = 0.02729$, $\alpha_p(\text{Si}) = 0.025$, $\alpha_s(\text{F}) = 0.1049$, $\alpha_p(\text{F}) = 0.0826$, $\alpha_s(\text{Cl}) = 0.05048$, $\alpha_p(\text{Cl}) = 0.05087$, $\alpha_s(\text{Br}) = 0.0469096$, and $\alpha_p(\text{Br}) = 0.0465342$. The final contraction schemes were as follows: H(5s1p/3s1p), C(10s6p1d/5s3p1d), Si(13s9p1d/7s5p1d), F(10s6p1d/5s3p1d), Cl(13s9p1d/7s5p1d), and Br(15s12p6d/9s7p3d). To benchmark the DFT methods employed in this study, more rigorous CCSD(T)^{30,31} computations were carried out on SiH₃ with the five silicon core orbitals frozen and no restriction on excitation into the virtual space. These SiH₃ CCSD(T) computations were carried out in conjunction with the Dunning–Woon aug-cc-pVQZ^{32,33} basis set. For SiH₂F, with the aug-cc-pCVTZ³⁴ basis set, we correlated all electrons and placed no restriction on the virtual orbitals. All CCSD(T) computations were performed using the ACESII³⁵ suite of programs.

The three types of neutral anion energy differences reported in this article were evaluated as differences of total energies: the adiabatic electron affinity

$$\text{AEA} = E(\text{optimized neutral}) - E(\text{optimized anion})$$

the vertical electron affinity

$$\text{VEA} = E(\text{optimized neutral}) - E(\text{anion at optimized neutral geometry})$$

and the vertical detachment energy of the anion

$$\text{VDE} = E(\text{neutral at optimized anion geometry}) - E(\text{optimized anion})$$

III. Results and Discussion

Selected equilibrium geometries are presented in Figures 1–3, with the remainder shown in the Supporting Information. The ground electronic states of all the neutral silyl radicals in this work are doublets, whereas the corresponding anions are closed-shell singlets. The computed values of the AEA, VEA, and VDE were corrected for the zero-point vibrational energy and are listed in Tables 1–3 respectively; experimental values of the AEA, where available, are given in Table 1.

III.A. SiH₃ and SiH₃[−]. Equilibrium geometries for the ²A₁ ground state of SiH₃ and the corresponding ¹A₁ ground state of SiH₃[−] are shown in Figure 1; for SiH₃, the spin densities clearly indicate that its unshared electron is located primarily on the Si atom. All the functionals employed in this study overshoot both the Si–H bond distance and the H–Si–H bond angle of SiH₃ compared to the proposed experimental values,³⁶ 1.46 ± 0.01 Å and 110 ± 2°. A similar tendency was reported by Pak

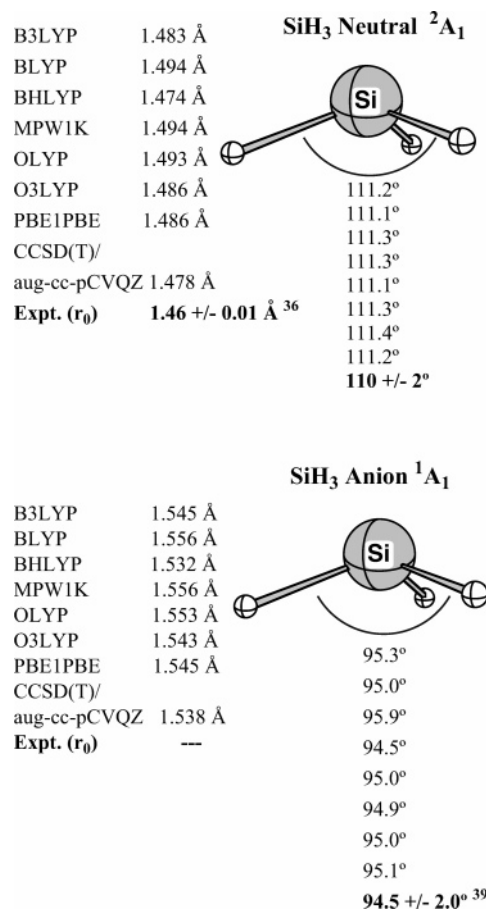


Figure 1. Equilibrium geometries for the ²A₁ state of the radical SiH₃ and the ¹A₁ state of the SiH₃[−] anion.

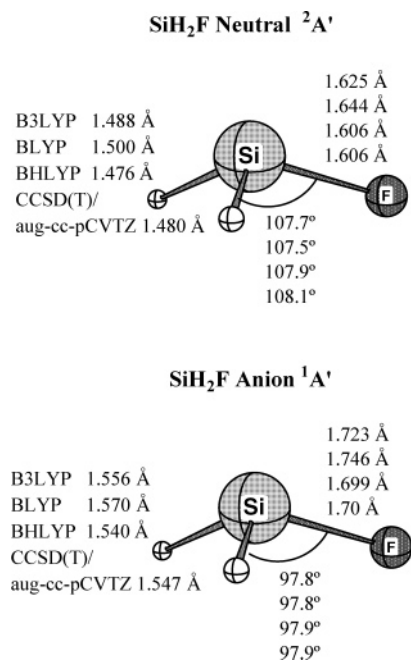


Figure 2. Equilibrium geometries for the ²A' state of the radical SiH₂F and the ¹A' state of the SiH₂F[−] anion.

et al.³⁷ in a prior investigation of unsubstituted silicon hydrides. It should be noted, however, that these DFT results are in good agreement with our high-level CCSD(T)/aug-cc-pVQZ results, where the computed Si–H bond length, 1.481 Å, and H–Si–H bond angle, 111.3°, are also greater than the analogous

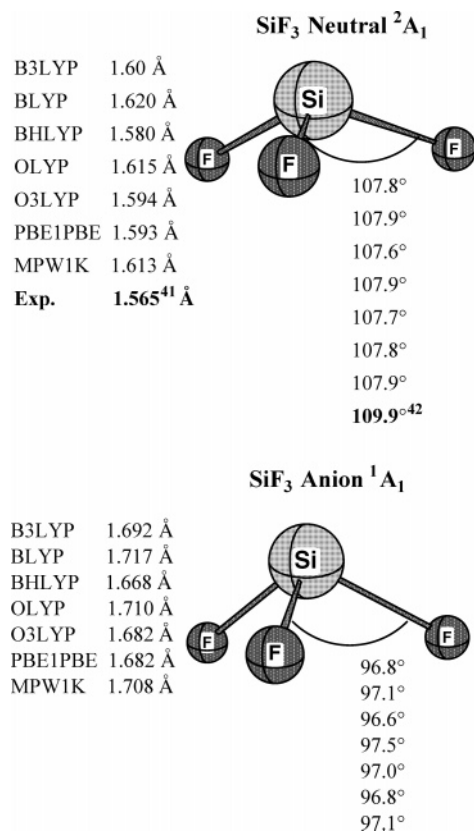


Figure 3. Equilibrium geometries for the ²A₁ state of the SiF₃ radical and the ¹A₁ state of the SiF₃⁻ anion.

experimental values. Of the functionals employed in this study, BHLYP predicts the shortest Si–H distance for SiH₃, 1.474 Å, although it is still 0.006 Å longer than the experimental *r*₀ value. The newer OLYP, O3LYP, PBE1PBE, and MPW1K functionals all give geometrical parameters for SiH₃ that are similar to those from the conventional BLYP, B3LYP, and BHLYP functionals (see Figure 1).

There are considerable changes in both the experimental and computed structures of SiH₃ upon the addition of an electron. In particular, the H–Si–H bond angle decreases sharply from its nearly tetrahedral value of ~111° in SiH₃ to a more steeply pyramidal value of ~95° for SiH₃⁻, see Figure 1. This decrease is primarily a result of doubly occupying a nonbonding orbital that has substantial *s*-character;³⁸ indeed, NBO analysis suggests that this lone-pair Kohn–Sham orbital has ~50% *s*-character. Our optimizations also predict that the additional electron in SiH₃⁻ results in an increase in the length of the Si–H bonds by ~0.06 Å. This increase helps alleviate the effects of hydrogen repulsion that would otherwise increase as the H–Si–H angle decreases.

The adiabatic electron affinities computed for SiH₃ in this study differ by as much as 0.23 eV, see Table 1; the MPW1K (1.44 eV), B3LYP (1.45 eV), and BLYP (1.37 eV) functionals provide the best agreement with the experimental value, 1.406 ± 0.014 eV, as determined by photoelectron spectroscopy.³⁹ It should be noted that our more rigorous CCSD(T)/aug-cc-pVQZ computations predict the AEA value to be 1.38 eV.

III.B. SiH₂CH₃, SiH₂SiH₃, SiH₂F, SiH₂Cl, SiH₂Br, and Their Anions. We first consider the monosubstituted silyl radicals SiH₂CH₃, SiH₂F, SiH₂Cl, SiH₂Br, SiH₂SiH₃, and their associated anions. Structures from our DFT optimizations of SiH₂F are shown in Figure 2, while those of the other monosubstituted species are given in the Supporting Information.

Unfortunately, there are no experimental structural data available for any of these monosubstituted species. Our computations predict that replacement of one of the hydrogen atoms with a fluorine atom results in only an ~0.01 Å elongation of the two remaining Si–H distances, while the H–Si–H angle increases by ~3.5°, see Figure 2. The predicted Si–H and Si–X (X = F, Cl, Br) distances using the BHLYP functional are consistently shorter than the corresponding distances using either the BLYP or the B3LYP functionals; see the Supporting Information.

Examination of Table 1 for these monosubstituted silyl radicals shows distinct trends in the computed values of the AEA compared to that of SiH₃. For example, for the BLYP functional, –CH₃ substitution decreases the AEA by 0.28 eV, whereas –SiH₃, –F, –Cl, and –Br substitution increase it by 0.40, 0.17, 0.57, and 0.68 eV, respectively. For the anions in this group of compounds photodetachment data³⁸ is available only for SiH₂CH₃⁻, and the resulting experimental value of the AEA, 1.01 ± 0.03 eV, is nearly 0.40 ± 0.04 eV (28 ± 3%) lower than the corresponding experimental value for SiH₃. For comparison, we note that the computed values of AEA for SiH₂–CH₃ are roughly 0.3 eV (~22%) less than the corresponding theoretical values for SiH₃.³⁹ Although the average difference between the magnitude of the experimental and computed values of SiH₂CH₃ is relatively large, 0.08 ± 0.05 eV, the O3LYP and PBE1PBE functionals perform very well, with predicted values of 1.02 and 1.04 eV, respectively. It should be noted that the BHLYP functional consistently predicts the lowest value of AEA for all the monosubstituted silyl radicals in this study; this is also the case for multiple substitutions, see Table 1. Our experience with this functional on a large set of compounds⁸ suggests that it typically provides lower bounds to experimental AEA results.

Although no experimental electron affinity data is available for SiH₂F or SiH₂Cl, our theoretical AEA values obtained from various DFT methods are in reasonable accord with those computed for these compounds by Rodriguez and Hopkinson⁴⁰ using isogyric reactions at the MP4/6-311++G(d,p)//HF/6-311++G(d,p) and MP4/6-311++G(2df,p)//HF/6-311++G(d,p) computational levels. Our BHLYP value of 1.40 eV for SiH₂F and our BLYP value of 1.94 eV for SiH₂Cl are in particularly good agreement with their results. We also carried out computations of the AEA for SiH₂F at the CCSD(T)/aug-cc-pCVTZ level, and the predicted value of 1.46 eV is also in good agreement with the predicted value of Rodriguez and Hopkinson.⁴⁰ Our BHLYP computed Si–H bond length, 1.476 Å, and bond angle, 107.9°, are in excellent agreement with the corresponding CCSD(T) values of 1.480 Å and 108.1°, respectively.

III.C. SiH(CH₃)₂, SiH(SiH₃)₂, SiHF₂, SiHCl₂, SiHBr₂, and Their Anions. Optimized structures of all the disubstituted radicals and their anions may be found in the Supporting Information. In accord with our observations from the previous section, the trend in electron affinities for these radicals follows a similar pattern. For example, for the BLYP functional, substitution with a second –CH₃ group lowers the predicted values of the AEA compared to that of SiH₂(CH₃) by 0.17 eV. Substitution with a second –SiH₃, –F, –Cl, or –Br group increase the AEA values by 0.28, 0.32, 0.59, or 0.62 eV, respectively, when compared to that of the corresponding monosubstituted radical. For the disubstituted silyl radicals, the only experimental electron affinity from the literature is for SiH(CH₃)₂; interestingly, the value of AEA, 0.91 ± 0.02⁵ eV, is only 0.1 eV (~9.9%) smaller than the experimental value for SiH₂(CH₃). Thus, the observed decrease in the AEA value in

TABLE 1: DFT/DZP++ and Experimental Adiabatic Electron Affinities (AEA) in eV

	B3LYP	BLYP	BHLYP	OLYP	O3LYP	PBE1PBE	MP1WK	exptl
SiH ₃ ^a	1.45	1.37	1.22	1.24	1.33	1.34	1.44	1.406 ± 0.014 ^b
SiH ₂ CH ₃	1.16	1.09	0.93	0.95	1.02	1.04	1.14	1.01 ± 0.03 ^c
SiH ₂ F ^{d,e}	1.63	1.54	1.40					
SiH ₂ Cl ^f	2.04	1.94	1.85					
SiH ₂ Br	2.16	2.05	1.97					
SiH ₂ (SiH ₃)	1.85	1.77	1.61					
SiH(CH ₃) ₂	0.98	0.92	0.74	0.78	0.83	0.84	0.95	0.91 ± 0.02 ^g
SiHF ₂ ^h	1.95	1.86	1.73					
SiHCl ₂	2.64	2.53	2.43					
SiHBr ₂	2.80	2.67	2.63					
SiH(SiH ₃) ₂	2.13	2.05	1.89					
Si(CH ₃) ₃	0.90	0.86	0.65	0.72	0.76	0.74	0.86	0.824 ± 0.02 ^e
SiF ₃	2.73	2.66	2.50	2.43	2.51	2.52	2.63	2.41 ± 0.22 ⁱ
SiCl ₃	3.31	3.21	3.10					
SiBr ₃	3.41	3.29	3.24					
Si(SiH ₃) ₃	2.34	2.25	2.10					
SiFCIBr	3.24	3.13	3.03					

^a CCSD(T)/aug-cc-pVQZ 1.38 eV. ^b Ref 39. ^c Ref 38. ^d CCSD(T)/aug-cc-pCVTZ 1.46 eV. ^e MP4/6-311++G(d,p)//HF/6-311G(d,p) 1.39 eV (ref 40). ^f MP4/6-311++G(d,p)//HF/6-311G(d,p) 1.94 eV (ref 40). ^g Ref 5. ^h MP4/6-311++G(d,p)//HF/6-311G(d,p) 1.82 eV (ref 40). ⁱ Ref 45.

TABLE 2: DFT/DZP++ Vertical Electron Affinities (VEA) in eV

	B3LYP	BLYP	BHLYP	OLYP	O3LYP	PBE1PBE	MPW1K
SiH ₃	0.96	0.88	0.72	0.76	0.85	0.84	0.94
SiH ₂ CH ₃	0.61	0.55	0.37	0.42	0.49	0.48	0.58
SiH ₂ F	1.02	0.96	0.77				
SiH ₂ Cl	1.29	1.21	1.06				
SiH ₂ Br	1.39	1.31	1.17				
SiH ₂ (SiH ₃)	1.20	1.11	0.96				
SiH(CH ₃) ₂	0.37	0.34	0.11	0.22	0.25	0.22	0.35
SiHF ₂	1.27	1.23	1.01				
SiHCl ₂	1.71	1.65	1.47				
SiHBr ₂	1.90	1.84	1.67				
SiH(SiH ₃) ₂	1.40	1.31	1.17				
Si(CH ₃) ₃	0.31	0.32	0.01	0.21	0.22	0.13	0.29
SiF ₃	2.69	1.90	1.63	1.70	1.75	1.69	1.86
SiCl ₃	2.26	2.24	1.98				
SiBr ₃	2.49	2.46	2.22				
Si(SiH ₃) ₃	1.58	1.49	1.35				
SiFCIBr	2.24	2.22	1.95				

TABLE 3: DFT/DZP++ Vertical Detachment Energies (VDE) in eV

	B3LYP	BLYP	BHLYP	OLYP	O3LYP	PBE1PBE	MPW1K
SiH ₃	1.91	1.81	1.68	1.68	1.79	1.83	1.91
SiH ₂ CH ₃	1.65	1.56	1.43	1.42	1.51	1.54	1.63
SiH ₂ F	2.13	2.02	1.93				
SiH ₂ Cl	2.73	2.58	2.57				
SiH ₂ Br	2.84	2.67	2.71				
SiH ₂ (SiH ₃)	2.46	2.37	2.23				
SiH(CH ₃) ₂	1.48	1.39	1.26	1.24	1.32	1.35	1.44
SiHF ₂	2.47	2.33	2.29				
SiHCl ₂	3.40	3.21	3.29				
SiHBr ₂	3.52	3.29	3.45				
SiH(SiH ₃) ₂	2.82	2.71	2.57				
SiF ₃	3.45	3.32	3.27	3.09	3.20	3.26	3.31
SiCl ₃	4.15	3.96	4.03				
SiBr ₃	4.13	3.91	4.19				
Si(SiH ₃) ₃	3.05	2.92	2.81				
Si(CH ₃) ₃	1.37	1.29	1.16	1.13	1.20	1.24	1.32
SiFCIBr	4.04	3.85	3.94				

going from SiH₂(CH₃) to SiH(CH₃)₂ is nearly a factor of 2 smaller than that predicted from our DFT results. The magnitude of the computed values for SiH(CH₃)₂ differ from experiment by 0.08 ± 0.05 eV on average, which is similar to what we found for SiH₂(CH₃). The BLYP functional gives an AEA value of 0.92 eV for SiH(CH₃)₂, which is indistinguishable from the experimental result (see Table 1). For SiHF₂, the computed value of the AEA at the BLYP, 1.86 eV, is very close to the MP4/

6-311++G(d,p)//HF/6-311++G(d,p) result of Rodriguez and Hopkinson,⁴⁰ 1.82 eV.

III.D. Si(CH₃)₃, Si(SiH₃)₃, SiF₃, SiCl₃, SiBr₃, and Their Anions. Equilibrium geometries for the ²A₁ ground electronic state of the SiF₃ neutral radical and the corresponding ¹A₁ closed-shell anion are presented in Figure 3; geometries of the other trisubstituted species are given in the Supporting Information. Tanimoto and Saito⁴¹ have reported a microwave structure

TABLE 4: Changes in the BLYP/DZP++ AEAs, Using SiH₃ as a Point of Reference

Effect on the AEA of SiH ₃ from Each Substituent	
SiH ₃	0.00 ^a
SiH ₂ CH ₃	-0.28
SiH(CH ₃) ₂	-0.4
Si(CH ₃) ₃	-0.51
SiH ₂ SiH ₃	0.40
SiH(SiH ₃) ₂	0.68
Si(SiH ₃) ₃	0.88
SiH ₂ F	0.17
SiHF ₂	0.49
SiF ₃	1.29
SiH ₂ Cl	0.57
SiHCl ₂	1.16
SiCl ₃	1.84
SiH ₂ Br	0.68
SiHBr ₂	1.30
SiBr ₃	1.92

^a The actual predicted AEA of SiH₃ is 1.37 eV.

for SiF₃; the Si–F bond length is deduced to be 1.565 Å, when the F–Si–F bond angles are assumed to be 109.9°, from the IR matrix isolation study of Milligan et al.⁴² As noted previously, the functionals in this study typically overshoot the Si–X (X = H, F) bond lengths when compared to experiment, and for SiF₃ the predicted Si–F distances are from 0.015 to 0.055 Å longer than the infrared/microwave result. On the other hand, the F–Si–F bond angle, in contrast to the H–Si–H angle in SiH₃, is less than the experiment by 2.0–2.3°. Again, the BLYP functional predicts a structure that is closest in agreement to experiment, see Figure 3. Our DFT optimizations predict a significant change in the structure of the neutral SiF₃ radical upon the addition of an electron. The F–Si–F angle decreases dramatically, on average by 10.8°, and this is accompanied by an increase in the Si–F bond length of ~0.09 Å; this increase in bond length reduces the effects of fluorine substitution in this compound.

In accord with an observation made by Wetzel et al.,³⁸ we note that methyl substitution is not additive in the sense that, the large decrease, 0.4 eV (experimental), 0.3 eV (theoretical), in the AEA in going from SiH₃ to SiH₂(CH₃) is followed by a much smaller reduction, 0.2 eV (experimental), 0.1 eV (theoretical), in going from SiH₂(CH₃) to SiH(CH₃)₂. Replacing the final H atom with a methyl group results in an even smaller reduction in the AEA, see Table 4. This trend was also noted by Frenking et al. for AX₃⁺ and AH₂X⁺ systems (A = C, Si, Ge, Sn, Pb; X = F, Cl, Br, I).⁴³

It is evident from Tables 1–3 that for mono-, di-, and trisubstituted silyl radicals the electron affinity consistently depends on the substituents in the order CH₃ < H < F < SiH₃ < Cl < Br. The rate of increase in electron affinity for the halo and silyl substituents is also shown in Table 4. These results are consistent with those found for the electron affinities of methyl radicals and silylenes,^{9,10} where a similar inverse relationship between the electronegativity of a substituent and its corresponding electron affinity was observed.

An explanation of this trend for the halogenated silyl radicals relies on the fact that fluorine is a weaker π -donor than the other halogens. Frenking and co-workers⁴³ showed that for halogens, the π -donor ability increases with F < Cl < Br < I for both cations and neutral molecules, and this increase in electron density causes silicon to become a much stronger electron attractor. Interestingly enough, the same effect can be seen among the relationship between methyl and silyl substituents

due in part to the electropositive nature of silicon^{38,44} as well as the π -back-donation described by Frenking et al.⁴³ This rationalization can also be extended to the inverse relationship of electron affinities and electronegativities observed previously^{9,10} for the divalent CX₂ and SiX₂ molecules (X = CH₃, SiH₃, F, Cl, and Br).

III.E. Performance of Newer Functionals. Considering the wealth of new functionals^{15–19} being reported in the literature and in view of our ongoing research effort to find efficient, accurate methods to predict electron affinities, we have considered four of the more recently developed functionals to assess their ability to predict EAs. Unfortunately, only limited experimental data is available to compare with for the silyl radicals discussed in this work.

For the parent SiH₃ radical, the MPW1K functional performs better than the other newer OLYP, O3LYP, and PBE1PBE functionals, as well as the exhaustively tested BLYP, B3LYP, and BLYP functionals, providing the closest agreement to the experimental value.³⁹ The MPW1K functional has an interesting property: for the five silyl radicals where experimental data is available, the predicted AEAs from this functional are all greater than the observed values. If this property persists on a more comprehensive database, the MPW1K functional may prove useful in generating upper bounds on AEA values. The B3LYP functional also appears to have the same general property.⁸

Although the O3LYP functional performs better than all the other functionals in predicting the AEA of SiH₂CH₃ (see Table 1), its performance is erratic in that in some cases it gives AEA values greater than experiment, while in other cases it gives values less than experiment. Interestingly, AEA values obtained from O3LYP are consistently lower than those from B3LYP; similarly AEA(OLYP) < AEA(BLYP).

For both SiH(CH₃)₂ and Si(CH₃)₃ the best agreement with experiment is obtained using the well-tested BLYP functional. However, the predicted AEA values from the OLYP, O3LYP, and PBE1PBE functionals are only ~0.10 to 0.15 eV too low. Furthermore, for all four of the newer functionals AEA(Si(CH₃)₃) < AEA(SiH(CH₃)₂), which is consistent with experiment.³⁸

In the case of SiF₃, the OLYP functional gives an AEA value that is closest to experiment. However, the uncertainty associated with the experimental electron affinity of SiF₃, ±0.22⁴⁵ eV, is quite large, making it difficult to determine which functional gives results closest to experiment. Indeed, only two of our computed AEA values actually fall outside the reported error. It should also be noted that the BLYP and OLYP functionals give values of AEA that are below the experimental values for SiH₃, SiH₂(CH₃), SiH(CH₃)₂, and Si(CH₃)₃ but above that for SiF₃. This could indicate that the experimental value of 2.41 eV for SiF₃ may be somewhat low. In view of these limitations in the experimental result for SiF₃, particularly the large error bars (which are nearly an order of magnitude greater than that associated with the other experimental EAs), we have excluded the AEA of SiF₃ in our assessment of average absolute errors. For the ensemble of silyl radicals studied in this work, DZP++ BLYP provides the most accurate computations of electron affinities—the average absolute error is only about four-hundredths of an eV. BLYP has been shown previously to provide excellent results for a large set of compounds.⁸ The average errors in increasing order are as follows: BLYP (0.04₁) < O3LYP (0.05₈) < MPW1K (0.06₀) < PBE1PBE (0.06₃) < B3LYP (0.08₅) < OLYP (0.11₅) < BLYP (0.15₃). Although the results presented in this article are less than exhaustive, the performance of the O3LYP, MPW1K, and PBE1PBE function-

als in computing AEAs is encouraging. Further testing on a more diverse collection of compounds is clearly required to assess the performance these functionals in computing EAs.

IV. Conclusions

Substituted silyl radicals were examined in this research to characterize the effect that alkyl and halo substituents have on the magnitude of electron affinities, as well as to provide an ensemble of molecules to gauge the accuracy of several newer pure (OLYP) and hybrid (O3LYP, MPW1K, PBE1PBE) density functionals for future electron affinity studies. An inverse relationship between electronegativities and electron affinities is found for the silyl radicals in this study, which is consistent with analogous studies of methyl radicals, carbenes, and silylenes.^{9,10} This, we assert, is a consequence of back-donation of electron density via π -bonds from the heavier halo substituents. This effect stabilizes the neutral structures and increases their attraction for electrons.⁴³ Thus, substituents with weak back-donating character, e.g., fluorine, have lower than expected values of the AEA.

Interestingly, the rather old BLYP method^{11,14} performed the best of all the functionals in this study in computing AEAs for those silyl radicals where reliable experimental results are available for comparison. However, the more recently developed O3LYP, PBE1PBE, and MPW1K hybrids show promise in this regard, although extensive testing on a more diverse collection of molecules is clearly required. In general, the performance of DFT methods for predicting molecular electron affinities of a variety of structures^{9,10} is encouraging, particularly in light of the great difficulty⁸ in determining experimental EAs to within ± 0.1 eV.

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Supporting Information Available: BLYP and B3LYP structures. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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